

# Phase ordering after a deep quench: Glauber dynamics for the Ising and hard-core gas on trees

F. Martinelli, Univ. Roma 3 (joint work with P. Caputo)

# Introduction

- $G = (V, E)$  a countable infinite graph (bdd degree)
- $P_t \Leftrightarrow$  Markov semigroup of a Glauber dynamics on  $G$  with multiple reversible extremal Gibbs measures  $\mu_1, \mu_2, \dots$
- On finite balls the  $\mu_i$ 's usually selected by boundary conditions (e.g. all  $+/-$  for Ising model).
- With b.c.  $\Leftrightarrow \mu_i$ , even if the Glauber starts from a very atypical (for  $\mu_i$ ) point  $\eta$ ,  
 $\delta_\eta P_t(\cdot) \approx \mu_i(\cdot)$  for  $t > T_{\text{mix}}$ .

- **An important open problem:** is the dynamics on  $G$  able to select a specific  $\mu_i$  as  $t \rightarrow \infty$  if the initial configuration is sampled from a suitable distribution  $\nu$  ? If so how fast ?
- Interesting case:  $\nu$  highly disordered (product Bernoulli or “high-temperature” Gibbs measure).
- Finite setting formulation  $(G_n, P_t, \mu_i)$ : for suitable  $\nu$  are the marginals of  $\delta_\eta P_t$  very close to  $\mu_j \neq \mu_i$  for  $1 \ll t \ll T_{\text{mix}}$  and “most”  $\eta$  w.r.t.  $\nu$  ?

→ **Example:** the Ising model on  $G = \mathbb{Z}^d$ ,  $d \geq 2$ .  
If  $\beta > \beta_c$  and  $\eta = \text{all “+”} \Rightarrow \delta_\eta P_t \rightarrow \mu^+$  while if  
 $\eta = \text{all “-”} \delta_\eta P_t \rightarrow \mu^-$ .

→ **Conjecture (Liggett):** if  $\{\eta_x\}_{x \in \mathbb{Z}^d}$  i.i.d. with  
 $p \equiv \nu(\eta_0 = +1) > \frac{1}{2}$

$$\nu P_t \rightarrow \mu^+ \quad \text{as } t \rightarrow \infty$$

(convergence = convergence of marginals of finite subsets of  $V$ .)

→ **Major obstacle:** multiple Gibbs measures  $\Rightarrow$   
poor analysis of mixing times on finite balls.

- ➔ On  $\mathbb{Z}^d$  or hexagonal lattice several results for Glauber dynamics for the Ising model at zero temperature ( $\beta = \infty$ ) [Howard-Newman], [Newman-Stein], [Camia-De Santis-Newman], [Fontes-Schonmann-Sidoravicius].
- ➔ For our problem:  $p \approx 1 \Rightarrow$  stretched exp. fixation [Fontes-Sidoravicius-Schonmann]:

$$\int d\nu(\eta) \mathbb{P}(\sigma_t^\eta(0) = -1) \leq e^{-t^{1/d}}$$

Related to competition between two bootstrap percolation processes.

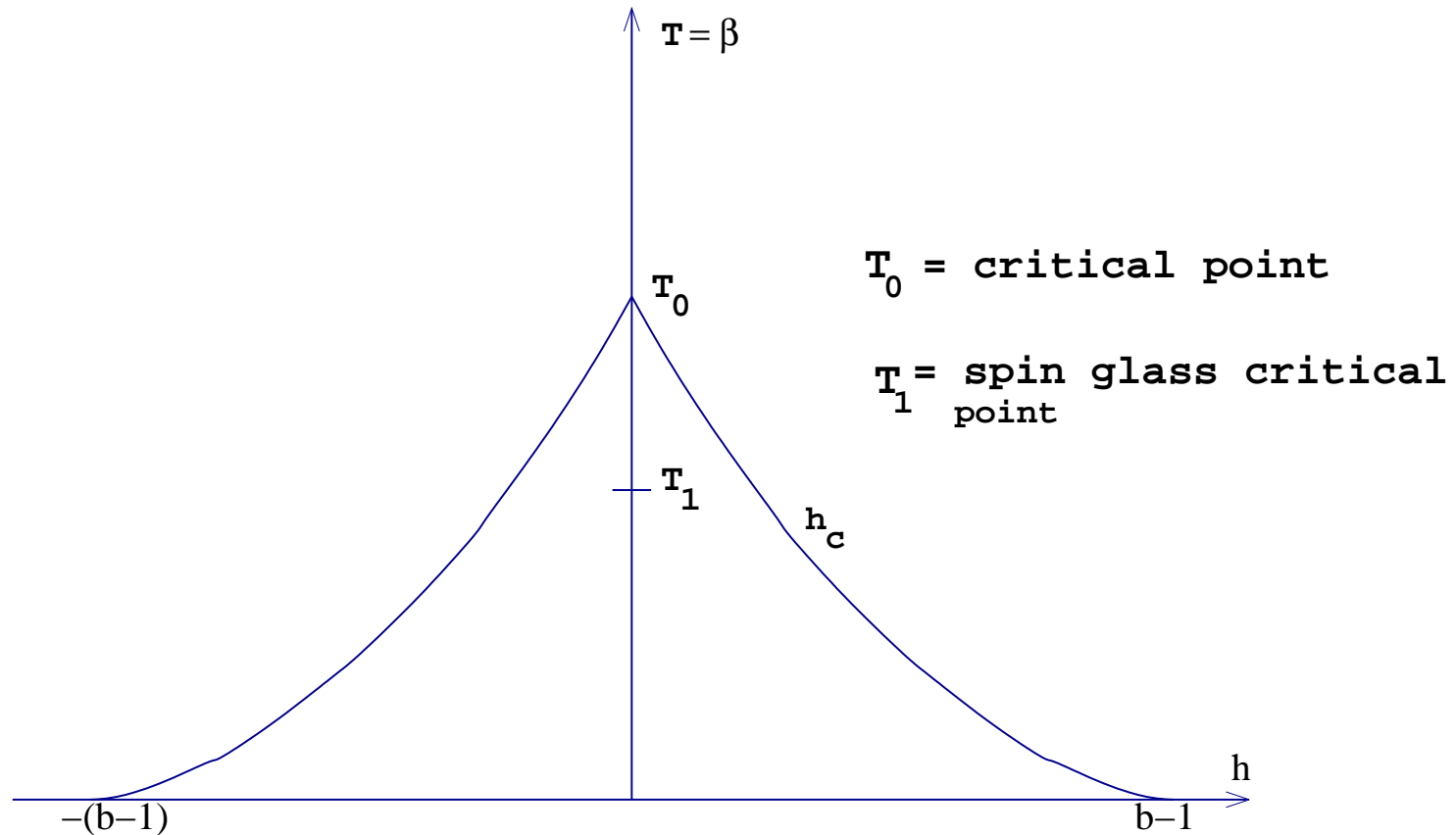
# The Ising model on trees

- $T \subset \mathbb{T}^b$ , the  $b$ -ary tree;  $E$  edge set of  $\mathbb{T}^b$ .  
 $\partial T$  consists of the children of its leaves.
- $\Omega = \{-1, +1\}^{T \cup \partial T}$ ,  
 $\Omega^\tau = \{\sigma \in \Omega; \sigma|_{\partial T} = \tau|_{\partial T}\}$
- The Gibbs measure ( $\sigma \in \Omega^\tau$ ):

$$\mu_T^\tau(\sigma) \propto \exp \left[ \beta \left( \sum_{x,y \in E} \sigma_x \sigma_y + h \sum_{x \in T} \sigma_x \right) \right],$$

- $\lim_{T \uparrow \mathbb{T}^b} \mu_T^+ = \mu^+$ ;

# The Phase Diagram



➔  $\mu^+ \neq \mu^-$  if  $\beta > \beta_0$  and  $|h| \leq h_c(\beta)$ .

# Glauber dynamics

On  $\Omega = \Omega_T$  consider the following Markov process  $t \rightarrow \sigma_t^\eta$  started from  $\eta$ :

- ➔ attach independent Poisson clocks of rate one, one to each vertex  $x$ ;
- ➔ when the clock at  $x$  rings replace  $\sigma_x$  with a new value out of the conditional Gibbs measure on  $x$  conditioned on the current values of its neighbors.
- ➔  $(P_t f)(\eta) := \mathbb{E}(f(\sigma_t^\eta))$  for local  $f$ .



# Main results: preliminaries

**Definition 1.**  $\Omega_\alpha$  will denote the set of starting configurations  $\eta \in \Omega$  such that for any  $x \in \mathbb{T}^b$  there exists a time  $t_0 = t_0(\eta, x) < \infty$  such that for all  $t \geq t_0$

$$| \mathbb{E}(\sigma_t^\eta(x)) - \mu^+(\sigma(x)) | \leq \exp(-t^\alpha).$$

- ➔ If  $\eta \in \Omega_\alpha$  then  $\delta_\eta P_t \rightarrow \mu^+$  weakly (convergence of marginals) as  $t \rightarrow \infty$ .
- ➔ Call  $\mathbb{P}_p$  the product Bernoulli measure with  $p := \mathbb{P}_p(\eta_x = +1)$ .

# Main result

**Theorem 2.** *In any of the following three cases*

$\exists \alpha > 0$  s.t.  $\nu(\Omega_\alpha) = 1$  for any  $\nu \geq \mathbb{P}_p$ :

i)  $p$  large enough, any  $\beta$  but  $h > -h_c(\beta) + \epsilon(p)$   
TM  
.

ii)  $p > \frac{1}{2}$ ,  $b$  large enough, any  $\beta$  large enough and  
 $h = 0$ .

iii)  $p > 0$  (even tiny),  $b$  large enough,  $\beta$  large,  
 $h = h_c(\beta)$ .

(!) For  $\beta$  large and  $h = h_c(\beta)$ , starting from all  $-1$  leads you to the minus phase.

# Remarks

- ➔ The event  $\Omega_\alpha$  is increasing:  
 $\eta \leq \eta'$  and  $\eta \in \Omega_\alpha \Rightarrow \eta' \in \Omega_\alpha$ .

$$\mathbb{E}(\sigma_t^\eta(x)) \leq \mathbb{E}(\sigma_t^{\eta'}(x)) \leq \mathbb{E}(\sigma_t^+(x))$$

Using the fast mixing for  $\mu^+$  proved by [Sinclair-Weitz-M.]  $\Rightarrow$

$$0 \leq \mathbb{E}(\sigma_t^+(x)) - \mu^+(\sigma(x)) \leq Ce^{-mt}$$

- ➔ Enough to consider  $\nu = \mathbb{P}_p$ .

# Remarks

- ➔ Not surprising (because of fast mixing)  
 $\mu^+(\Omega_{\alpha=1}) = 1 \Rightarrow \text{any } \nu \geq \mu^+ \text{ has } \nu(\Omega_1) = 1.$
- ➔ For any  $(\beta, h)$   $\mathbb{P}_p \geq \mu^+$  if  $p \approx 1$  (depending on  $\beta, h$ ).
- ➔ Hard part is to prove the result for *all*  $\beta$  large enough.
- ➔ Even if  $c_{\text{sob}}(\mu^+) > 0$  (Weitz talk), methods based on decay (in time  $t$ ) of relative entropy of  $\nu P_t$  w.r.t.  $\mu^+$  work only if  $\text{Ent}\left(\frac{d\nu}{d\mu^+}\right) < \infty.$

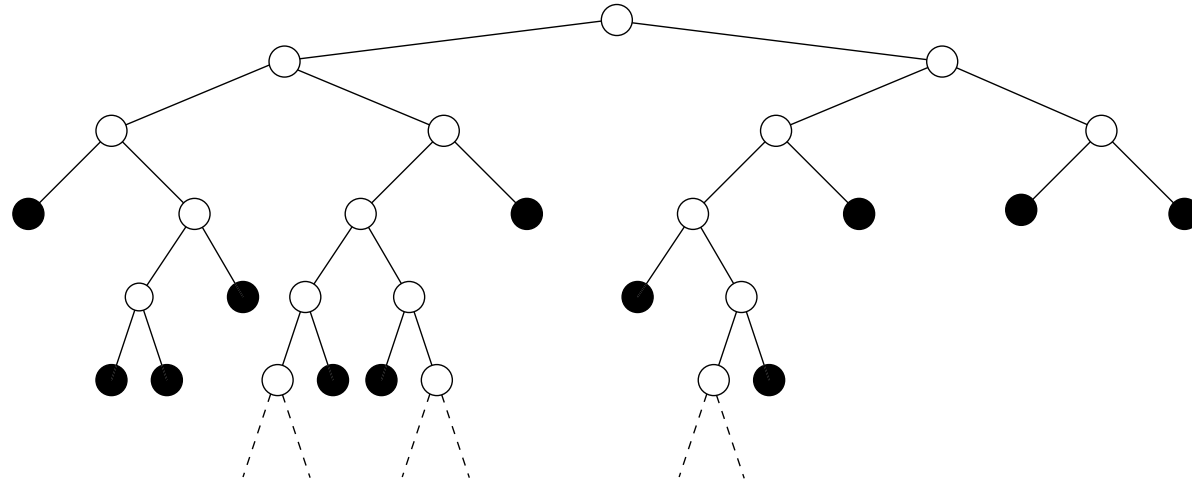
# Main strategy for the proof

Need a *lower bound* on (restrict to root  $r$ ):

$$\mathbb{E}(\sigma_t^\eta(r)) - \mu^+(\sigma(r))$$

- Fix a length scale  $\ell = t^\alpha$  and freeze (in time) all the  $\eta(x) = -1$   $\ell$  levels below the root;
- Call  $\omega$  the configuration of “obstacles” (= frozen vertices);

$$\begin{aligned} \mathbb{E}(\sigma_t^\eta(r)) - \mu^+(\sigma_r) &\geq \text{(by monotonicity)} \\ [\mathbb{E}(\sigma_{t,\omega}^\eta(r)) - \mu_\omega^+(\sigma_r)] - [\mu^+(\sigma_r) - \mu_\omega^+(\sigma_r)]. \end{aligned}$$

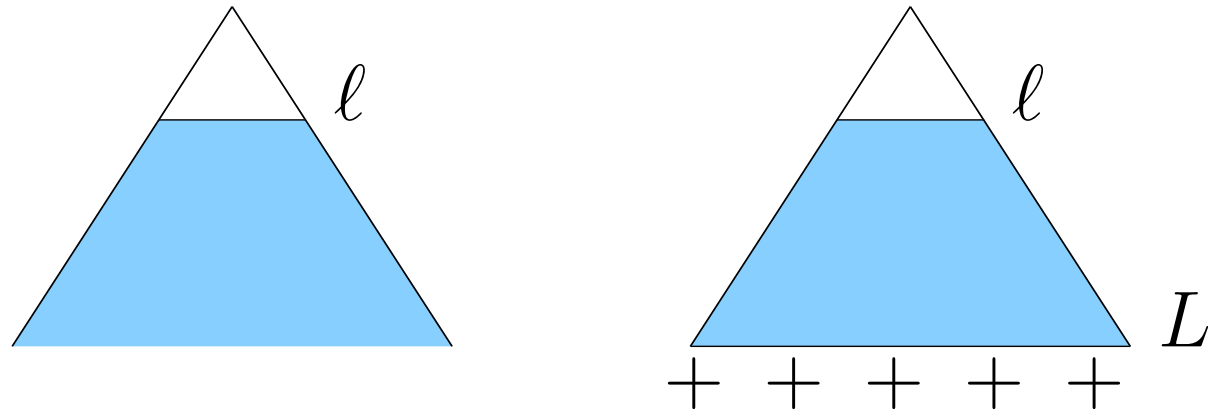


Free vertices ( $\circ$ ) and obstacles ( $\bullet$ ) in a given realization  $\omega$ .

- Notice that  $\eta = +1$  on the vertices of random tree ( $\circ$  vertices)  $\mathbb{T}_\omega^b$ ;
- Analogies with the original problem with  $\eta = +1$  but on a *random* tree.

# From infinite to finite tree

Fix another length scale  $L = \ell^\gamma$ ,  $\gamma > 1$ ;



Random obstacles below level  $\ell$ : Infinite tree (left) and finite tree with  $+$  boundary condition below level  $L = \ell^\gamma$ ,  $\gamma > 1$  (right).

- ➔ Again monotonicity  $\Rightarrow \mu_\omega^{+, (L, +)} \geq \mu_\omega^+$  and  $\mathbb{E}(\sigma_{t, \omega}^\eta(r)) \leq \mathbb{E}(\sigma_{t, \omega}^{\eta, (L, +)}(r))$ ;

# Conclusion

$$\mathbb{E}(\sigma_t^\eta(r)) - \mu^+(\sigma_r) \geq \quad (\text{previous step})$$

$$[\mathbb{E}(\sigma_{t,\omega}^\eta(r)) - \mu_\omega^+(\sigma_r)] - [\mu^+(\sigma_r) - \mu_\omega^+(\sigma_r)] \\ \geq$$

$$[\mathbb{E}(\sigma_{t,\omega}^{\eta,(L,+)}(r) - \mu_\omega^{+,(L,+)}(\sigma_r)] \quad \text{(A)}$$

$$- [\mathbb{E}(\sigma_{t,\omega}^{\eta,(L,+)}(r) - \mathbb{E}(\sigma_{t,\omega}^\eta(r))] \quad \text{(B)}$$

$$- [\mu^+(\sigma_r) - \mu_\omega^+(\sigma_r)] \quad \text{(C)}$$



# The three steps

- (A)  $\rightarrow$   $[\mathbb{E}(\sigma_{t,\omega}^{\eta,(L,+)}(r) - \mu_{\omega}^{+,(L,+)}(\sigma_r))]$  needs mixing times bounds for a random tree with obstacles (the frozen vertices  $-1$ );
- $\rightarrow$  As in [Sinclair-Weitz-M] but now log-sob  $c_{\text{sob}}(\mu_{\omega}^{+,(L,+)})$  shrinks as  $L^{-c}$  with large  $\mathbb{P}_p$ -probability;
- (B)  $[\mathbb{E}(\sigma_{t,\omega}^{\eta,(L,+)}(r) - \mathbb{E}(\sigma_{t,\omega}^{\eta}(r)))]$  requires a true dynamical analysis (coupling);
- (C)  $[\mu^+(\sigma_r) - \mu_{\omega}^+(\sigma_r)]$  easier. Recursive analysis on random tree.

# Why does it work ?

1. Attractivity (some monotonicity w.r.t. the partial order  $\sigma \leq \eta$  iff  $\sigma_x \leq \eta_x$  for all  $x$ ).  
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2. Rigidity of the extremal phases  $\mu^\pm$ . If  $h > -h_c(\beta)$  adding a small density of obstacles (extra  $-1$ ) below  $\ell$  does not change significantly  $\mu^+$  above  $\ell$ .

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3. Rapid mixing (exp fast) of the Glauber dynamics w.r.t  $\mu^+$ .

# Starting from $p = 1/2$

- What happens when  $\beta > 0$ ,  $h = 0$  and  $p = \frac{1}{2}$  ?
- At  $\beta = +\infty$  infinite (+) and (-) chains form in finite time and freeze [Howard];
- Recall the spin glass critical point

$$e^{-2\beta_1} = \frac{\sqrt{b}-1}{\sqrt{b}+1} \cdot \text{TM}$$

**Conjecture 0.** *If  $\beta < \beta_1$  and  $p = 1/2$  then  $\nu_{1/2}$  a.s.  $P_t^\eta \rightarrow \mu^0$ , the Gibbs measure obtained as  $\lim_{\Lambda \rightarrow \mathbb{T}^b} \mu_\Lambda$  with no boundary conditions .*

# Supporting the conjecture

- ➔ Random (w.r.t.  $\mathbb{P}_{p=\frac{1}{2}}$ ) boundary conditions  $\Leftrightarrow$  free b.c.
- ➔ The random variable  $P_t f(\eta) := \mathbb{E}(f(\sigma_t^\eta(r)))$  becomes a sure r.v. :

$$\mathbb{P}_{\frac{1}{2}} \left( \eta; |P_t f(\eta) - \mathbb{E}_{\frac{1}{2}}(P_t f)| \geq e^{-ct} \right) \leq e^{-c_f e^{ct}}.$$

Notice that for odd (w.r.t.  $\sigma \rightarrow -\sigma$ )  $f$ 's  
 $\mathbb{E}_{\frac{1}{2}}(P_t f) = \mu^0(f) = 0$ .

- ➔ For  $\beta < \beta_1$  the free measure  $\mu^0$  is extremal;

# Proof of concentration

- ➔ Proof via standard concentration bounds for  $\mathbb{P}_{\frac{1}{2}}$ .

$$\mathbb{P}_{\frac{1}{2}}(\eta; |F(\eta)| \geq r) \leq e^{-\frac{r^2}{2}}$$

for any mean zero function  $F$  with unitary Lipschitz norm

$$\|F\|_{\text{Lip}}^2 := \sum_{x \in \mathbb{T}^b} \|F(\eta^x) - F(\eta)\|_{\infty}^2.$$

- ➔ Need bounds on Lipschitz norm of  $P_t f$ .

- Coupling + weighted Hamming distance technique of [Berger, Kenyon, Mossel, Peres] + [Peres-Winkler]  $\Rightarrow$

$$\|P_t f(\eta^x) - P_t f(\eta)\|_\infty \leq C_f e^{-ct} \lambda^{d(x)}$$

$$b\lambda^2 < 1 \text{ if } \beta < \beta_1 \Rightarrow \|P_t f\|_{\text{Lip}} \leq C'_f e^{-ct}$$

- Concentration leads to stability of  $\mathbb{E}_{\frac{1}{2}}(P_t f)$  w.r.t. perturbations of  $\mathbb{P}_{\frac{1}{2}}$  with moderate exponential growth as  $\ell \rightarrow \infty$  of relative entropy between marginals on first  $\ell$  levels.