

8/15/05

Seiberg-Witten equations &  
dynamics of vector fields in dim 3  
Clifford Taubes

$M = 3$ -dim'l oriented manifold  
 $\Omega = \text{volume}$

(#1)  $v \lrcorner \Omega = \omega_v$      $d\omega_v$     Vector field dynamics

$\phi_t : M \rightarrow M$  diffeomorphisms  
 $\frac{d\phi}{dt} = v(\phi_t)$

Contact 1-form  $a \in \Omega^1(M)$

$$\Omega = a \wedge da > 0$$

$$v \lrcorner da = 0$$

$$v \lrcorner a = 1$$

$$\omega_v = da$$

Thm:  $v$  has at least 1 closed orbit.

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Stable Hamiltonian (joint w/ Hutchings)

$$a \quad \omega_v \quad a \wedge \omega_v = -\Omega$$

$$da = f\omega_v$$

In this case,  $v$  has a closed orbit

except if  $M = T^2$   
 $\downarrow$   
 $S^1$

What can be said about uniquely ergodic vector fields?

Say that a vector field is exact if  $\omega_v = db$

asymptotic linking number:  

$$S_v = \int_M b \wedge \omega_v$$

Thm: If  $v$  is uniquely ergodic, then  $S_v = 0$ .

Example:

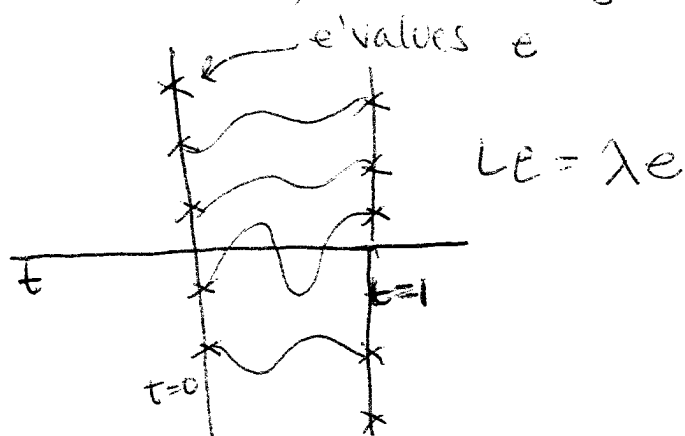
Horocycle flow on unit tangent bundle to the surface

$$\omega_v = db$$

$$b \wedge db = 0$$

## #2 Spectral Dynamics

$\mathbb{H}$  - Hilbert space  
 $L: \mathbb{H} \rightarrow \mathbb{H}$ , self-adjoint



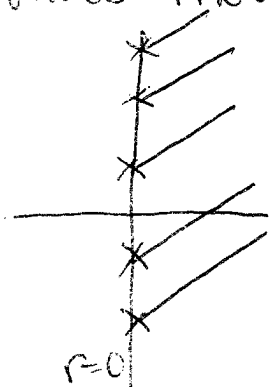
Spectral flow  $\{L_t\}_{t \in [0,1]}$

Simple example:

$$L_r = i \frac{\partial}{\partial \theta} + r$$

$$\{e^{in\theta}\}_{n \in \mathbb{Z}} \quad \lambda = -n + r$$

eigenvalues move linearly



$$\int_0^1 f = r + \mathcal{O}(1)$$

spectral flow

Fix metric on  $M^3$ ,  $\Omega$   
 $|v|=1$  and  $|wv|=1$

$$F \supseteq U(2) \supseteq U(1)$$

$$\begin{array}{ccc} \downarrow & & \downarrow \\ Fr & \supseteq & SO(3) \end{array}$$

$$Spin^c \iff H^2(M; \mathbb{Z})$$

$$\downarrow \\ M$$

$$F \times_{U(2)} \mathbb{C}^2 = \mathbb{S}$$

$$TM \xrightarrow{\phi} \text{End}(\mathbb{S})$$

$$cl(b)cl(c) = -\langle b, c \rangle - cl(* (b \wedge c))$$

clifford multiplication

Dirac operator

$$C^\infty(M, \mathbb{S}) \xrightarrow{\nabla_A} C^\infty(M; \mathbb{S} \otimes T^*M) \xrightarrow{cl} C^\infty(\mathbb{S}; M)$$

DA-irb

$$v \lrcorner \Omega = wv = db$$

Spectral flow is independent of  $b$   
 if  $c_1(\det \mathbb{S}) = \text{torsion}$

Thm:  $\delta f = \frac{-1}{8\pi^2} r^2 \int b \wedge \omega_V + O(r^{2-\delta})$   
 $\delta > \frac{1}{32}$

$$b \wedge db = 0$$

### (#3) Seiberg-Witten dynamics

Spin<sup>c</sup> structure  $A$ , a connection on  $\det(S)$   
 $\psi$ , a section of  $S$

- $D_A \psi = 0$

- $*F_A = r(\psi^+ \lrcorner \psi - r * \omega_V)$

$$\langle c, \eta^+ \lrcorner \psi \rangle = \eta^+ \text{cl}(c) \psi$$

Thm:

Suppose  $\{(A_n, \psi_n), r_n\}_{n=1,2,\dots}$

with

$$r_n \rightarrow \infty$$

- $E(A_n) = i \int * \omega_V \wedge F_{A_n} < \epsilon < \infty$

- $\sup_M |r_n| > \delta > 0$

then  $V$  has a closed orbit.

Suppose  $E(A_n) \rightarrow \infty$

Define  $f \rightarrow \frac{\int f(1-|\gamma_n|^2)}{E(A_n)}$

Then this converges as  $n \rightarrow \infty$  to an invariant measure which is nontrivial if  $A \neq 0$

Remarks:

If  $c_1(\det S)$  is torsion, then there exists lots of  $\{(A_n, \gamma_n), r_n\}$  obeying SW-equations.

SW-Floer homology.

Cycles come from solutions to SW equations

Irreducible soln:  $C^\infty(M, S')$  on solutions

$$A \rightarrow A - 2u^{-1}du \quad \gamma \rightarrow u\gamma$$

stabilizer = 1,  $\gamma \neq 0$

Reducible soln:  $S^1, \gamma \equiv 0$

$$A = \int_{\text{flat}} -irb' \quad db = \omega v$$

Thm (Kronheimer, Mrowka)

There are infinitely many non-trivial  
classes with degrees  $\rightarrow -\infty$

# Day 3 Kinetics of v. field dynamics

## Part I: continuous space dynamics:

- $M =$  closed oriented  $k$ -d mfd.
- $V =$  vector field

$V$  generates 1 param. family of diffeos.  
 $x \mapsto \phi_t(x)$

$$\frac{d\phi}{dt} = V(\phi_t) \quad \phi_0 = x$$

Interest is in behavior of

$$\phi_t(x) \text{ as } t \rightarrow \infty.$$

- Are there closed orbits? ( $\phi_T(x) = x$  for some  $T > 0$  &  $x$  non-trivially)
- Are there inv. measures?  $d_\mu^* S^2 \sim S^2$  the same



3 PARTS:

PART 1: Vector field dynamics:

1.  $M =$  compact oriented 3 manifold  
 $\alpha =$  contact 1 form on  $M$  so  $\alpha \wedge d\alpha > 0$   
 $v =$  Reeb vector field,  $v \lrcorner d\alpha = 0$   $v \lrcorner \alpha = 1$

Of interest are closed orbits of  $v$ .

Thm (W. Thurston) - Any such  $v$  has a closed orbit.

What if  $v$  is not of "contact type"?

Stable Hamiltonian system

$\alpha, f, \omega$                        $\alpha$  is 1-form

$f$  is function

$\omega$  is 2-form

$d\alpha = f \omega$

$\alpha \lrcorner \omega > 0$

So  $f$  can be zero, pos, neg.

$v \lrcorner \omega = 0$   
 $v \lrcorner \alpha = 1$

Thm (Mike Hutchings & ... )  $v$  has a closed orbit except ~~possibly~~ maybe

if  $M \cong T^2$ .

$\downarrow$   
ST

How about your general work?

• Solve just  $\exists \int \omega$  Volume preserving  $V$ -  
Hamiltonian. ... flows both to  
K. Einsmann ... closed orbits

• If it has no zero, orbit, what can  
we say about it?

• If in real dynamical systems, Every  
trajectory is dense?  
Gottschalk conj.  $\exists$  none on  $\mathbb{Z}^d$  ??

• If  $v$  is volume preserving then  
 ~~$L_v \Omega = 0$~~   $L_v \Omega = 0$  so  $\int \Omega$  gives  
invariant measure. A vector field  
is said to be "volume preserving"  
if  $\int \Omega$  is an invariant measure  
via  $L_v \Omega = 0$ .  
It follows that  $\int \Omega$  is constant.

What can be said about  $\int \Omega$ ?  
... ..

• Only a few are known to exist:

• eg. Int. tangent vector of surface of genus  $g \geq 1$ . Homocycle flow is uniquely ergodic.

See that  $v$  is uniquely ergodic if  $v \perp \Omega = db$

(e.g. homocycle flow is exact.)

Define  $\Delta_v = \int_M b_1(v \perp \Omega)$  (Asymptotic winding #)

It is not independent of anti-div.  $b$

• Thm (C.T): If  $v$  is uniquely ergodic then  $\Delta_v = 0$ .

\* Winding # (winding # of any curve  $\gamma$  is  $\int \gamma \perp \Omega$ )

\* Any flow on a  $\mathbb{Q}$ -homology  $\mathbb{S}^1$  is exact.

\* Homocycle flow is exact & in fact

$$v \perp \Omega = db \quad \& \quad b \text{ and } b' = 0$$

So  $v$  is tangent to a foliation

(leaves are gen. by  $v$  & geodesic flow.)

PART 2: Vector field

• Spectral dynamics

$H = L^2(\text{circle})$  space

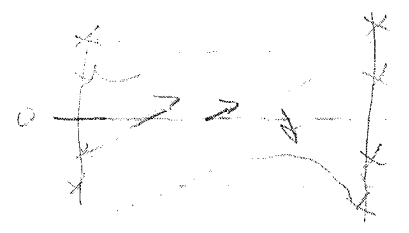
$L$ : Self-adjoint operator on  $H$  with domain  $D(L)$  dense in  $H$ .

$L$  has discrete spectrum in  $\mathbb{R}$  w/ no accumulation pt. finite mult.



•  $\{L_x\}_{x \in \mathbb{R}}$  are family of such op. with  $0 \notin \text{spec}(L_0), \text{Spec}(L_1)$

Spectral flow



$\parallel$  were  $\neq$  det.

Ex:  $L_n = \left[ i \frac{\partial}{\partial \theta} + n \right]$  on  $L^2(\mathbb{R}/2\pi\mathbb{Z})$

Spec  $L_n = \{n + r\}$

SD  $\dots = \mathbb{R} + \mathbb{Z} (!)$

Exact fibration on  $M^3$  Put Riem. metric with vol. form =  $\Omega$

$$F = U(2)$$

$\downarrow$        $\downarrow$

$$F_0 = SO(2)$$

$\downarrow$

$\mathbb{R}$

Sphere bundle

$$\Leftrightarrow \mathbb{H}^3 / U(1)$$

$$F \rightarrow \mathbb{R} \rightarrow \mathbb{S}$$

Sphere bundle  
Module for Clifford mult

$$T^*M \xrightarrow{cl} \text{End}(\mathbb{S})$$

$$cl(\partial_t) = cl(\partial_t)$$

$$cl(\partial_i) = -\langle \partial_i, \cdot \rangle + \langle \cdot, \partial_i \rangle$$

$$\det \mathbb{S} = F \times_{\det} \mathbb{O} \quad (= \frac{1}{2} \mathbb{S})$$

Conn  $A$  on  $\det \mathbb{S} \Rightarrow$  Conn. deriv on  $\text{Spin}^c$

$$D_A = cl(\nabla_A) = \text{S.ad}_\rho \text{ on } \mathbb{P}(N; \mathbb{S})$$

Fix  $\mathbb{P} = \text{conn on } \det \mathbb{S}, \quad \mathbb{D} = 1\text{-Conn on } \mathbb{H}$

$v \rightarrow D_{\mathbb{P} \oplus \mathbb{D}}$  ~~we~~ look at special flow

$$\text{eg. } d\mathbb{D} = \omega_v = v \lrcorner \Omega$$

Note: if  $cl(\det \mathbb{S}) = \text{tr conn}$ , then  $\mathbb{S}$  gives unit rep. of group of anti-herm.

Ans: 
$$\text{sf } \vec{E} \cdot \vec{r} = \frac{1}{4\pi^2} r^2 \int b \omega \omega_V + \mathcal{O}(r^2 \omega^3)$$
  
 (Asymptotic expansion #1).

Part 2:  $\omega \rightarrow \infty$

• Write the vector potential  $\vec{A}(\vec{r}, t)$  in terms of  $\vec{r}$  and  $\omega$  when  $\omega \rightarrow \infty$ .  $\text{div } \vec{A} = 0$ .  $\text{curl } \vec{A} = \vec{E}$ .

ex. for sublight  $\text{sf } \vec{r} \cdot \vec{r}^2 \text{ vol.}$

Part 3: Non-linear domain of LSW of vector field.

Solve in  $\mathcal{N}$  eqn.

$A = \text{curl } \vec{\Psi}$

$\Psi = \text{curl } \vec{\Phi}$

$$\nabla \times \vec{A} = \vec{r} (\nabla^2 \Psi - i a)$$

$$\nabla^2 \Psi = 0$$

$$\langle \vec{b}, \nabla^2 \vec{\Psi} \rangle = \nabla^2 \langle \vec{b}, \vec{\Psi} \rangle$$

Interest is in how solutions vary as  $r \rightarrow \infty$ .

(Are there any? of what sort? What happens as  $r \rightarrow \infty$ ?)



• Critical pts of  $\alpha \Leftrightarrow$  solutions of SW

• Floor inv =  $\mathbb{Z}$  module generated by  $\psi$  in part of equiv. classes of solutions

$$A_0 \psi \rightarrow A - 2\psi^2, \psi \rightarrow \psi^4 \quad \psi: M \rightarrow S^1$$

Mod.  $\Rightarrow \psi \neq 0$  then  $\psi^2 = 0$  then  $\psi^4 = 0$

Data in equiv. classes and in  $H^2(M; \mathbb{Z})$

$\mathbb{Z}$ -equiv

• Then  $\langle \text{ker} - M \rangle \in \mathbb{Z}$  ...  $\langle \text{SW} \rangle \in \mathbb{Z}$

### Discussion about Floor homology

What is point? It quantize solutions to an equation!

If  $\text{ker} \neq \emptyset$ , there must be some solutions.

(Evens fact. on compact world has at least  $\sum b_i$  crit. pts.)

• Floor inv. Floor inv. do take account of ...





Thus,

$$\mu_n^2 = \frac{\int (U - \mu_n)^2}{\int (U - \mu_n)} \quad \text{has } \mu_n^2 = 1$$

$$\mu_n^2 \rightarrow 0$$

$$\text{But } \int \delta^2 = \mu \neq 0$$

These expressions are proved in  
d. or a with it.

Long story.